Ion conduction in vanadium-substituted $LiSn_2P_3O_{12}$ electrolyte nanomaterials

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Abstract $\text{LiSn}_2\text{P}_3 - {}_y\text{V}_y\text{O}_{12}$ powders with y = 0.2, 0.4,0.6, and 0.8 are prepared by mechanochemical milling method. The pellets of the compounds are heat treated at temperatures between 700 to 1,000 °C for sintering period of 8 h. X-ray diffraction analysis indicates that all samples consist of rhombohedral crystalline LiSn₂P₃O₁₂ phase. Energy dispersive X-ray analysis confirmed that V^{5+} has been successfully substituted into LiSn₂P₃O₁₂ crystalline phase. The conductivities of the pellets are determined using impedance spectroscopy. Impedance analysis shows enhancement in both bulk and grain boundary conductivities with increase in y. The enhancement in bulk conductivity is due to decrease in bulk activation energy reflecting an increase in ion mobility as a result of an increase in bottleneck size. Enhancement in grain boundary conductivity is attributed to increase in the number of conducting pathways due to an increase in crystallite homogeneity.

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Introduction

Polycrystalline materials with sodium super ionic conductor (NASICON) structure exhibit high ionic-conducting capabilities due to existence of three dimension, 3D, conduction pathways which enables easy migration of mobile ions [1-4]. LiSn₂(PO₄)₃ is one of the compounds that possess structure analogous to NASICON structure. The basic structure of $LiSn_2(PO_4)_3$ is rhombohedral. This basic structure consists of a three dimensional framework of two SnO₆ octahedra sharing corner oxygens of PO₄ tetrahedra forming $[Sn_2(PO_4)_3]^-$. This framework is covalently bonded making it highly stable with high melting point of above 1,650 °C [3]. In the $[Sn_2(PO_4)_3]^-$ skeleton, each SnO₆ octahedron is separated by three PO₄ tetrahedra and each of this PO₄ tetrahedra is connected to four SnO₆ octahedra [2]. This skeleton in turn forms chain like "ribbons" along the *c*-axis. The ribbons are joined together along the a- and b- axis by PO₄ tetrahedra. The corner sharing of oxygen ions between the SnO₆ octahedra and PO₄ tetrahedra results in an open framework structure with structural holes. These structural holes give some freedom for rotational motion to the structure. This flexibility allows substitution of various ions at Sn and P sites. $LiSn_2(PO_4)_3$ framework consists of two types of sites that can be occupied by lithium ions: (i) M1 sites located between SnO₆ octahedral are sites that have six coordination with oxygens forming LiO₆ octahedra and (ii) M₂ sites, located between the ribbons of PO₄ tetrahedra has a distorted eight coordination with oxygens. Both sites are arranged in an alternating form of a 3D array along conduction channels, also known as bottlenecks, which allow transportation of cations [1-6].

A few researchers have reported studies on $LiSn_2P_3O_{12}$ prepared via solid state reaction [7, 8]. However, their

LiSn₂P₃O₁₂ pellets showed low conductivities in the order of 10^{-10} S cm⁻¹ due to the presence of Teflon which was used as a binding agent to avoid pellets from breaking during characterization. Recently, the authors have reported stable pellets of LiSn₂P₃O₁₂ prepared by mechanochemical milling method without the use of any binder [9]. These pellets showed conductivity values in the order of 10^{-7} S cm⁻¹ which are higher than those reported in references [7, 8]. The focus of this work is to prepare vanadium-substituted V⁵⁺ into LiSn₂P₃O₁₂ to form compounds with a general formula of LiSn₂(PO₄)_{3 - y}(VO₄)_y with the hope of enhancing conductivity. The compositional effects on the conductivity of the compounds as well as their structure are investigated.

Experimental

Stoichiometric proportions of Li₂O (99%, Aldrich), SnO₂ (98%, Sigma-Aldrich), NH₄H₂PO₄ (98%, Sigma-Aldrich), and V_2O_5 (98%, Aldrich) were prepared according to the formula $\text{LiSn}_2\text{P}_3 - {}_y\text{V}_y\text{O}_{12}$ (y = 0.2, 0.4, 0.6, and 0.8) and heated at 700 °C for 2 h. This heating process was done to decompose the hydrogen phosphate and initiate its reaction with the carbonate during heating [8, 10]. After heating process, the mixtures were milled for 80 h using Fristch 7 planetary ball mill. Milled powders were then pelletized to form pellets with diameter of 13 mm using a hydraulic press at a pressure of 7 ton. The pellets were then sintered at four different temperatures, T from 700 to 1,000 °C for 8 h. It was observed that different y values exhibit optimum conductivity at different sintering temperatures which were at 700 and 600 °C for y = 0.2 and 0.4, respectively, while for y = 0.6 and 0.8, the optimum sintering temperature was 500 °C. For further analysis samples were prepared by sintering them at their optimized sintering temperatures and these samples are labeled as V2, V4, V6, and V8 for y of 0.2, 0.4, 0.6, and 0.8, respectively.

Impedances of samples were measured using Solartron SI 1260 impedance analyzer while X-ray diffraction (XRD) data of the samples were recorded using Bruker D8 Advance X-Ray Diffractometer. To study crystallite size, Scherer equation was applied [11];

$$D = \frac{k\lambda}{\beta\cos\theta} \tag{1}$$

where k is the Scherer constant value (0.94), λ is the wavelength of the source (1.5406 Å), β is FWHM (in radians), and θ is the Bragg angle (in radians). Quanta 200 FESEM scanning electron morphology (SEM) equipment interfaced with energy dispersive X-ray (EDX) instrument was used for microstructure and elemental composition analyses. Transference numbers were analyzed by Wagner's polarization

method using Solartron SI 1287 Electrochemical Interface in order to determine the type of charge carrier in the samples.

Results and discussion

The X-ray diffraction spectra for V2 to V8 are depicted in Fig. 1. The spectra clearly indicate the presence of rhombohedral LiSn₂(PO₄)₃ crystalline phase (label with miller indices of plane) in the samples. This phase has lattice parameters; a = 8.650 Å, b = 8.650 Å, and c = 21.487 Å. These parameters are same with the parameters obtained for LiSn₂P₃O₁₂ systems reported in [9] showing that the vanadium-substituted samples have the same crystalline phase as that of the $LiSn_2(PO_4)_3$ parent compound. SnO₂ impurity is also observed in all samples indicating the presence of unreacted SnO₂. However, no peak attributed to vanadiumrelated compounds is observed. Therefore, it can be inferred that vanadium has been successfully substituted into LiSn₂(PO₄)₃ crystalline matrix. Figure 2 shows the comparison of crystallite size for all samples. All samples are found to have crystallite sizes in the order of 10^{-9} m indicating that they are nano in size. The distribution of crystallite size in each sample is almost the same at all planes except for plane 300 especially for samples V2 and V4. This indicates that samples V6 and V8 have crystallites which are homogenous while those of V2 and V4 contain crystallites of smaller sizes in between larger crystallites.

Figure 3 shows micrograph images of the morphology for the cross-section surface of V2, V4, V6, and V8. V2 shows some agglomeration of grains of different sizes. In V4, the agglomerations are of smaller sizes and more homogeneous



Fig. 1 X-ray diffraction patterns for V2, V4, V6 and V8



Fig. 2 Variation of crystallite size with y (vanadium content) at different planes

compared to V2. V6 and V8 show crystallites of about the same size which is in agreement with the results of crystallite size analysis presented in Fig. 2. In V6, agglomerations of grains are also observed. In addition, needle-like structures are also found scattered between these agglomerations.

Figure 4 shows EDX spectra for V2, V4, V6, and V8. The details of the EDX spectra analysis are tabulated in Table 1. As lithium is not detectable by EDX due to its light atomic weight which does not permit the analysis of emitted radiations, the concept of charge neutrality is applied [12]. It is found that the ratios of the elements in the compounds are in good agreement with the stoichiometric ratio of the starting materials. This analysis confirms that vanadium has been successfully substituted at P sites in LiSn₂P₃O₁₂ matrix. This supports the XRD results discussed earlier. From the table, it can also be observed that the needle-like structures contain high percentage of vanadium. However, this compound is expected to exist in small amounts as it is not detectable by XRD. Structure of SnO₂ observed in XRD is not detected in SEM micrographs. This compound is believed to be present in other spots or parts in the samples which are not shown in Fig. 3.

A typical impedance plot for the studied compounds is displayed in Fig. 5. The figure shows two overlapping semicircles followed by a spike in the low frequency region. The high frequency semicircle is assigned to bulk response with its intercept at the *x*-axis assigned to bulk resistance, R_b , while the middle frequency semicircle is assigned to grain boundary response with its intercept at the *x*-axis corresponds



Fig. 3 SEM micrographs of V2, V4, V6, and V8



Fig. 4 EDX spectra at each points a, b, and c for all samples

Table 1 Atomic percentage of elements at each point for all samples

Sample	Composition	Stoichiometric ratio Sn:P:V
V2	Starting mixture	2.0:2.8:0.2
	Fig. 4 V2-a	2.00:2.79:0.19
	Fig. 4 V2-b	2.00:2.78:0.20
V4	Starting mixture	2.0:2.6:0.4
	Fig. 4 V4-a	2.00:2.56:0.41
	Fig. 4 V4-b	2.00:2.51:0.38
V6	Starting mixture	2.0:2.4:0.6
	Fig. 4 V6-a	2.00:2.35:0.61
	Fig. 4 V6-b	2.00:2.42:0.55
	Fig. 4 V6-c	1.00:1.28:2.94
V8	Starting mixture	2.0:2.2:0.8
	Fig. 4 V8-a	2.00:2.16:0.79
	Fig. 4 V8-b	1.00:1.22:2.90



Fig. 5 Impedance plot for V2

to grain boundary resistance, $R_{\rm gb}$. Bulk and grain boundary conductivities, $\sigma_{\rm b}$ and $\sigma_{\rm gb}$ are calculated using $R_{\rm b}$ and $R_{\rm gb}$ values determined from the impedance plots.

To confirm the conductivity obtained from the impedance plots, conductivity spectra (imaginary part of conductivity, σ'' vs. real part of conductivity, σ') are plotted and a typical spectrum is shown in Fig. 6. The conductivity plot consists of a semicircle and two dispersion curves at low σ' and high σ' regions, respectively. The intercept of the dispersion curves with *x*-axis at high and middle frequency gives the values of σ_b and σ_t (total conductivity), respectively. This method has been used by a few groups of researchers [13–16]. The σ_{gb} is obtained by subtracting σ_b from σ_t . The conductivities were then compared with the conductivity values obtained from the impedance plots and these values are found to be in good agreement with each other as listed in Table 2.



Fig. 6 Conductivity spectrum for V2

Table 2 Comparison of conductivity values determined from impedance and conductivity plots for all samples

Sample	$\sigma_{\rm b} \times 10^{-6} \ ({\rm S \ cm^{-1}})$		$\sigma_{\rm gb} imes 10^{-6} ({ m S \ cm^{-1}})$	
	Z plot	σ plot	Z plot	σ plot
V2	3.77 (±0.01)	3.76 (±0.03)	0.67 (±0.01)	0.65 (±0.01)
V4	4.81 (±0.03)	4.79 (±0.03)	1.51 (±0.02)	1.53 (±0.02)
V6	5.37 (±0.01)	5.40 (±0.03)	2.61 (±0.03)	2.61 (±0.01)
V8	5.70 (±0.02)	5.71 (±0.01)	2.71 (±0.01)	2.71 (±0.01)



Fig. 7 Conductivity variations with y for $LiSn_2P_{3-y}V_yO_{12}$ system

Figure 7 illustrates variations of bulk and grain boundary conductivities with vanadium content. As seen in the figure, both conductivities are observed to increase with



Fig. 8 Temperature dependence of conductivity for $LiSn_2P_{3-y}V_yO_{12}$ system

Table 3 Bulk and grain boundary activation energies for all samples

Sample	$E_{\rm a,b}~({\rm eV})$	$E_{a,gb}$ (eV)	
V2	0.17	0.30	
V4	0.16	0.29	
V6	0.14	0.30	
V8	0.12	0.29	

increase in vanadium content. V8 exhibits the highest conductivity with values of $\sigma_{\rm b}$ and $\sigma_{\rm gb}$ of 5.70 × 10⁻⁶ and 2.71 × 10⁻⁶ S cm⁻¹, respectively. This means that $\sigma_{\rm b}$ is increased by about four times while $\sigma_{\rm gb}$ shows an enhancement of one order of magnitude compared to $\sigma_{\rm b}$ and $\sigma_{\rm gb}$ of the LiSn₂(PO₄)₃ parent compound which are 1.65 × 10⁻⁶ and 1.05 × 10⁻⁷ S cm⁻¹, respectively [9].

Conductivity–temperature dependence of all samples is studied in the temperature range 353 to 463 K and the plots of log σ versus 1,000/*T* are shown in Fig. 8. It is observed that the bulk conductivity for all samples increases with temperature. All samples show linear log σ versus 1,000/ *T* plots with regression values of around 0.98 and 0.99 for

Fig. 9 Normalized current versus time for V2

both $\sigma_{\rm b}$ and $\sigma_{\rm gb}$, respectively. As such, it can be inferred that the conductivity is Arrhenian in nature. However, a sudden increase in $\sigma_{\rm b}$ is detected at around 170 °C for V4 and V6 which is believed to be associated with a change in crystalline phase [17]. The activation energy values, E_{a} were calculated from the Arrhenius slope and these are listed in Table 3. It is found that bulk activation energy, $E_{a,b}$ decreases with increase in vanadium content. V8 shows the lowest value of 0.12 eV. This value is smaller compared to E_a reported by Juarez et al. [7, 18] and Winand et al. [19] which are 0.34 and 0.55 eV, respectively. The decrease in activation energy reflects greater ion mobility which maybe due to an increase in bottleneck size in the structure [20, 21]. A similar observation has been reported by Chang et al. [21]. These authors reported that partial substitution of P^{5+} ions by V^{5+} ions in Li₁₃Al₀₃Ti₁₇(PO₄)₃ led to conductivity enhancement attributed to change in bottleneck size. In the system studied in the present work, $(PO_4)^{3-}$ tetrahedral is partially substituted by the larger $(VO_4)^{3-}$ forming larger bottlenecks leading to greater mobility of ions which in turn enhance conductivity [20, 21]. The increase in V content increases the bottleneck size leading to an increase in conductivity as shown in Fig. 7. In the case of grain boundary activation energy, $E_{a,gb}$, all samples show similar values indicating no significant effect of vanadium content on the $E_{a,gb}$. The increase in σ_{gb} with vanadium content maybe due to increase in the number of conducting pathways as a result of increase in contact surface area between grains as a result of increase in homogeneity of crystallites with increase in V content as presented in Fig. 2 [22, 23]. According to Kawai and Kuwano [24], σ_{gb} affects the ionic transportation in the bulk property. Thus, when σ_{gb} increase, the $\sigma_{\rm b}$ also increase as observed before.

The increase in the conductivity with vanadium content may also be attributed to increase in polarizability of cation–oxygen bond which formed the bottlenecks. According to Chowdari et al. [25] and Sugantha and Varadaraju [26], the larger the polarizability of cation-oxygen bond, the



easier the transport of ions in the bottleneck. Therefore, increase in vanadium increased the ion transportation which results in increase in conductivity.

Ionic transference number is the ratio of ionic conductivity to the total conductivity which is ionic and electronic. In this study, the ionic transference number is measured by Wagner's polarization method. Figure 9 presents a typical direct current polarization curve. Similar result obtains for all samples. The values of the transference numbers are 0.98, 0.98, 0.95, and 0.93 for V2, V4, V6, and V8, respectively. Since the transference numbers are close to unity, it can be inferred that the majority mobile charge carriers are ions and they are expected to be Li^+ . These values also indicate that there is a small electronic contribution associated with SnO₂ present in the samples as shown by XRD spectra in Fig. 1.

Conclusion

Partial substitution of V^{5+} in P^{5+} sites resulted in enhancement of bulk and grain boundary conductivities. Enhancement in bulk conductivity is due to increase in bottleneck size by partial substitution of vanadium which is reflected by a decreases in $E_{a,b}$. Increase in σ_{gb} with vanadium content is attributable to increase in number of conducting pathways as a result of increase in homogeneity of crystallites.

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